Current signatures and search for Pop. III stars in the Local Universe

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Abstract. Recent numerical studies argue that low-mass stars can be formed even in a zero-metallicity environment. These low-mass Population III (Pop. III) stars are thought to be still shining and able to be observed in the Local Universe. Most low-mass Pop. III stars are thought to be formed as secondary companions in binary systems. They can be escaped from their host mini-halos when their primary companions explode as supernovae. In this paper, we estimate the escape probability of the low-mass Pop. III stars from their host mini-halos. We find that $\sim$ 100 Pop. III stars are expected. We also compute spatial distribution of these escaped Pop. III survivors by means of a semi-analytic hierarchical chemical evolution model. Typically, they are distributed around $\sim$ 2Mpc away from the Milky Way but $5 - 35\%$ of the escaped stars fall into the Milky Way halo. These escaped Pop. III stars are possibly detected by very large scaled surveys being planned.

Key words. Stars: abundances – Stars: Population II – Galaxy: formation – early universe

1. Introduction

In spite of the longstanding effort, population III (Pop. III) stars, which are stars without metal, have not been discovered yet. Extremely metal-poor (EMP) stars in the Milky Way halo are searched by the large scaled surveys: HK survey, HES survey, and SDSS/SEGUE. These surveys found thousands of EMP stars with [Fe/H] < −3 but no star below [Fe/H] = −6. The observed most metal-poor objects, which are reffered to as hyper metal-poor (HMP) stars, have metallicity between [Fe/H] = −5 to −6 (HE1327-2326, [Fe/H] = −5.7, Frebel et al. 2005; HE0107-5240, [Fe/H] = −5.4, Christlieb et al. 2002).

On the other hand, from a theoretical perspective, the possibility of the formation of low-mass Pop. III stars has been pointed out. Recent numerical simulations argue that a primordial gas cloud fragments into multiple pieces, and form a Pop. III binary (Machida et al. 2008; Turk et al. 2009; Stacy et al. 2010) or a Pop. III star cluster (Clark et al. 2008, 2011; Greif et al. 2011). These results indicate that low-mass stars can be formed even in the pristine environment unaffected by prior star formation. Typical mass of the second generation Pop. III (Pop. III.2) stars, which are born from the gas ironized by the first generation stars, are thought to be smaller than the first generation Pop. III (Pop. III.1) stars because

of more efficient cooling (Uehara & Inutsuka 2000; O'Shea et al. 2005). If Pop. III stars with mass $m \lesssim 0.8\,M_\odot$ are formed, they are still shining by nuclear burning to date and can be observed in the Local Universe. We refer to these low-mass Pop. III stars which survive to date as Pop. III survivors.

In our previous studies, we investigated the change of the surface abundance by accretion of interstellar medium (ISM) on Pop. III survivors in the Milky Way halo (Komiya et al. 2009b, 2010). We found that the surfaces of Pop. III survivors are polluted to [Fe/H] $\sim$ -5. This result indicates that Pop. III survivors are observed as HMP stars when they are in the Milky Way halo. HMP stars are possibly Pop. III survivors but there are some other scenarios in which HMP stars are formed as chemically second generation stars (Umeda & Nomoto 2003; Iwamoto et al. 2005; Maynet et al. 2006).

Is it impossible to observe the unpolluted Pop. III survivors with zero surface metallicity in the Local Universe? In this paper, we propose that stars with $Z = 0$ can be found *outside* the Milky Way halo.

In the ΛCDM cosmology, Pop. III stars should be formed in the mini-halos with total mass $M_h \sim 10^6\,M_\odot$ (e.g. Yoshida et al. 2003). Low-mass Pop. III survivors have possibilities to escape from their host mini-halos because of the small gravitational potential of mini-halos in which Pop. III stars are formed.

Two channels to release Pop. III survivors can be considered. If a low-mass Pop. III star is formed as a secondary of a binary with a massive primary star, it will go out from the binary system when the primary explode as a SN. The other channel is the escape of a low-mass star through gravitational slingshot from a Pop. III star cluster. These stars can go away from mini-halo, and stay in intergalactic space. These "escaped Pop. III" stars are free from the surface pollution because it is difficult to accrete the metal enriched gas in the intergalactic space.

In this paper, we investigate the binary escape scenario. We evaluate the escape probability of Pop. III and EMP survivors from their host mini-halo and estimate the number and spatial distribution of them around the Milky
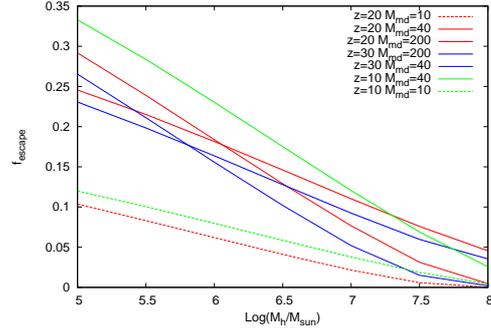
Fig. 1. Escape fraction of low-mass secondary stars of binaries as functions of mass of their host mini-halo. From a mini-halo of $m \sim 10^6\,M_\odot$, $\sim$ 20% of Pop. III survivors are escaped when the very high-mass IMF is adopted.

Way at $z = 0$. We also discuss the detection probability of the escaped Pop. III survivors.

2. Escape of Pop. III stars from their host mini-halo

When the primary star of a binary explode as a SN, the secondary star can be released from the binary system since the gravitational potential suddenly decreases. If the kinetic energy of the secondary star is larger than the sum of the gravitational potential of its host mini-halo and of the remnant object after the SN explosion,

$$\frac{1}{2}v_{orb}^2 - \frac{Gm_{rem}}{r_{bin}} + \Psi_{halo}(M_h, z) > 0, \quad (1)$$

the secondary star escape from the mini-halo, where v_{orb} and r_{bin} are the velocity of the secondary star and the separation of the binary, respectively. m_{rem} is the mass of the remnant (neutron star or black hole) of the primary star, and Ψ_{halo} is the gravitational potential of the mini-halo as a function of mini-halo mass and redshift. The terminal velocity of the escaped star becomes

$$v_{esc} = \left(v_{orb}^2 - \frac{2Gm_{rem}}{r_{bin}} + 2\Psi_{halo}(M_h, z)\right)^{1/2}. \quad (2)$$

For the gravitational potential of mini-halos, we adopt the NFW profile with the concentration parameter $c = 10$. It is assumed that

Pop. III stars are formed at the center of the mini-halos. m_{rem} is given as a function of initial mass of the primary following Woosley (2002). v_{orb} and r_{bin} are computed from mass of primary and secondary, and period of the binary assuming the circular orbit.

We estimate the escape fraction adopting the following parameter for the distribution of the primary mass, m_1, the mass ratio $q = m_2/m_1$, and the binary period.

We use the lognormal IMF. In our previous studies, we gave constraints on the IMF of EMP stars based on the statistics of carbon-enhanced EMP stars and the total number of observed EMP stars (Komiya et al. 2007, 2009a). As a result, we yield the lognormal IMF of medium mass $M_{md} = 10\,M_\odot$ and dispersion $\sigma = 0.4$ for EMP stars. Theoretical studies argue that the typical mass of Pop. III stars is larger than the EMP stars. We also adopt the very high mass IMF with $M_{md} = 40\,M_\odot$ or $M_{md} = 200\,M_\odot$ for Pop. III stars.

For the mass ratio distribution, we adopt two functions; a flat mass ratio distribution, $n(q) = 1$, as a fiducial one and a bottom heavy distribution, $n(q) = 2(1 - q)$. For the period distribution of binaries, we use the observational result of Duquennoy & Mayor (1991) as a fiducial one and adopt also one of Rastegaev (2010).

Secondaries of close binaries with period smaller than a few hundred yrs escape from mini-halos. Figure 1 shows the percentage of the escaped stars among Pop. III survivors as functions of mini-halo mass. When we assume the very high mass IMF with $M_{md} \geq 40\,M_\odot$, $\sim 20\%$ of the low-mass secondaries of Pop. III binaries escape from mini-halos with $10^6\,M_\odot$. In the case of $M_{md} = 10\,M_\odot$, the percentage becomes less than 10% because majority of primary companion of Pop. III survivors are intermediate massive stars. From more massive mini-halos with $10^8\,M_\odot$, the escape frequency is less than $\sim 5\%$ due to their deep gravitational potential.

3. Distribution of the escaped Pop. III stars in the Local Universe

In this section, we estimate the total number of the stars escaped from mini-halos as the building-blocks of the Milky Way and the spatial distribution of the escaped Pop. III or EMP stars. In our previous paper, we build a merger tree of the Milky Way halo using the extended Press-Schechter method by Somerville & Kolatt (1999), and follow the star formation history and the chemical evolution along the tree (Komiya et al. 2010; Komiya 2011; Komiya et al. 2014). We use the results of these previous studies.

We pick-up the Pop. III and EMP survivors formed in the hierarchical chemical evolution model, and set the binary parameters randomly following the distribution described above for each star. Stars which satisfy Eq. 1 escape from mini-halos.

We assume that the main halo of the merger tree has been stayed from the beginning at the position where the Milky Way locates now, and other mini-halos were accreted to the main halo, and compute the distance of the escaped Pop. III stars from the Milky Way.

For each escaped Pop. III star, we follow its radial orbit, $r_s(t)$, by integrating the equation of motion,

$$\frac{d^2 r_s}{dt^2} = -\frac{GM_{in}(r_s, t)}{r_s^2} + \frac{\Lambda c^2}{3} r_s + \frac{l_s^2}{r_s^3} \qquad (3)$$

where Λ is the cosmological term and $l_s \equiv r_{s,0} v_{ej} \sin\theta$ is the specific angular momentum of the star. We set the escape direction, θ, randomly. $M_{in}(r_s, t)$ is the total mass inside the radius r_s at time t and we assume that
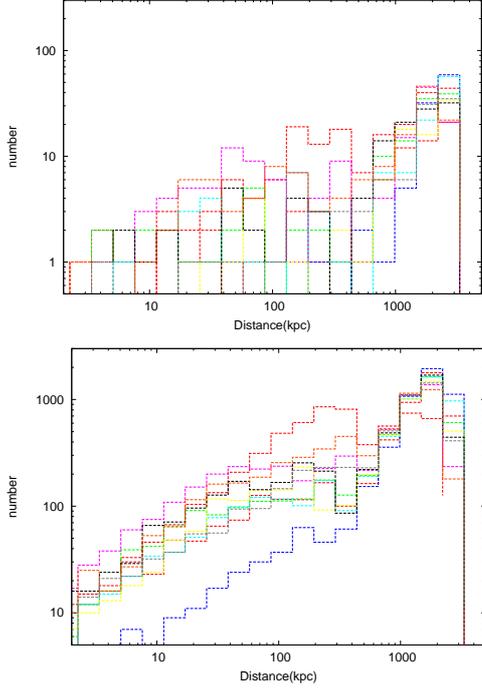
$$M(r_s, t) = \frac{4}{3}\pi r_s^3 \rho_{av}(t) + M_{main}(t) \qquad (4)$$

where M_{main} is mass of the main halo and ρ_{av} is the average density of the Universe. The initial distance and initial radial velocity are
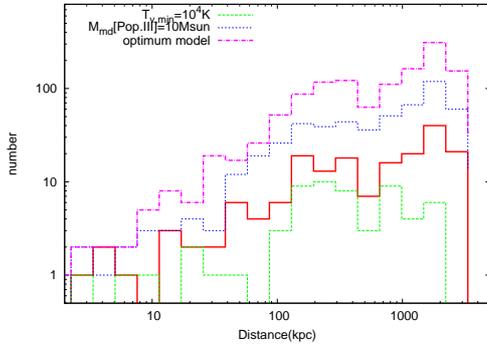
$$r_{s,0} = r_h \qquad (5)$$
$$v_{s,0} = v_{ej}\cos\theta + v_h \qquad (6)$$

where r_h and v_h are distance and radial velocity of the mini-halo in which the Pop. III star is formed.

Fig. 2. The predicted distribution of the escaped Pop. III survivors (top panel) and EMP survivors (bottom panel) around the Milky Way halo. We overplot the results of 10 computation runs for different merger trees

Fig. 3. Parameter dependence of the distribution of the escaped Pop.III stars.

The distance of the mini-halos is estimated by the following method. In the Press-Schechter formalism, "halo" is the collapsed spherical overdense region. In the extended Press-Schechter theory, "merger" of two halos

is described as a gravitational collapse of the region which contains both halos. Distance between the two halos which merge to form a halo with total mass, M_{tot}, at t_m is thought to be similar with a radius of the spherical shell of mass M_{tot} which collapse at t_m. Evolution of the spherical shell is simply described by the following equation,

$$\frac{d^2 r_h}{dt^2} = -\frac{GM_{tot}}{r_h^2} + \frac{\Lambda c^2}{3} r_h. \qquad (7)$$

We compute the distance, r_h, of a Pop. III forming mini-halo which merges with the main halo at t_m by integrating this equation with conditions of $r_h(t_m) = 0$ and $M_{tot} = M_{main}(t_m)$.

In the building blocks of the Milky Way, ~ 500 of Pop. III survivors are formed in total when we adopt $M_{md} = 40\,M_\odot$, $n(q) = 1$ and the period distribution of Duquennoy & Mayor (1991). More than 100 of them are escaped from their host mini-halos. They are formed at $z \sim 20$, and escaped after $\sim 10^7$ yrs with velocity $v_{ej} = 5 - 70$ km/s and $v_h = 50 - 250$ km/s. Top panel of Figure 2 shows the predicted distribution of the escaped Pop. III survivors. Typically, they are distributed around ~ 2 Mpc away from the Milky Way. $5 - 35\%$ of the escaped Pop. III survivors have accreted to the Milky Way and distributed around the virial radius of its dark matter halo.

We also compute the distribution of the escaped EMP survivors. More than 1000 EMP stars are escaped from mini-halos in our model. The distribution of them is similar to the escaped Pop. III stars, as shown in the bottom panel of Figure 2.

Figure 3 shows parameter dependence of the predicted distribution of the escaped Pop. III stars. When we adopt the optimum parameter set, $M_{md} = 10\,M_\odot$, $n(q) = 2(1-q)$, and the period distribution of Rastegaev (2010), the predicted number of the escaped Pop. III survivors becomes ~ 7 times larger than the fiducial case. On the other hand, when we assume that stars are formed only in halos with virial temperature higher than 10^4K, the number of the escaped Pop. III star reduce to $\sim 1/5$. For all cases, the escaped Pop. III stars are distributed around 2Mpc away from the Milky

Way. This is because the spatial distribution of these stars are dominated by the motion of their host mini-halos rather than the binary properties ($v_h > v_{ej}$) for most the escaped Pop. III stars.

The luminosity of Pop. III survivors at red giant branch or horizontal branch is $\sim 100 L_\odot$ of more. The observed magnitude of these stars outside the Milky Way is $\sim 26 - 28$ mag. There is a very wide ($\sim 1400 deg^2$) and deep ($i \lesssim 26$mag) survey using Hyper Suprime-Cam planned. Such a large scale survey possibly detects these escaped Pop. III or escaped EMP survivors.

4. Summary

We discuss the observation probability of the Pop. III stars in the Local Universe. We propose that the unpolluted Pop. III survivors can be observed outside the Milky Way halo. When the Pop. III binaries are formed, $\sim 20\%$ of low-mass secondary stars are thought to be escaped from their host mini-halos when their primaries explode as SNe. By means of the semi-analytic model for the star formation history of the Milky Way, we estimate the number and distribution of the escaped Pop. III survivors. We predict that $20 - 1000$ Pop. III survivors are distributed around ~ 2Mpc away from the Milky Way. $5 - 35\%$ of the escaped Pop. III stars fall to the Milky Way and is distributed around the virial radius of the dark halo. Thousands of EMP survivors may also be distributed at similar distance. The observational magnitude of these stars at ~ 2Mpc is $26 - 28$ mag when they are giant stars. There is a possibility to observe these stars by very wide and deep survey using such as Hyper Suprime-Cam.

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